

MAEDI-JSPS Sakura program

- A bilateral cooperative program between France and Japan
- To form sustained networks evolved from individual scientist exchanges **including young scientists**
- France {Marquet, Gelis, Iancu, Schoeffel, Soyez, Royon, Kohara, Petreska, Saimpert }
- Japan {Fujii, Hatta, Itakura, Nara, Sako, Yamazaki, Hagiwara, Taya, Zhou }





Update of heavy quark production in pA

H. Fujii
(U Tokyo, Komaba)

HF and K. Watanabe, NPA915(2013)1; 920(2013)78; 951(2016)45

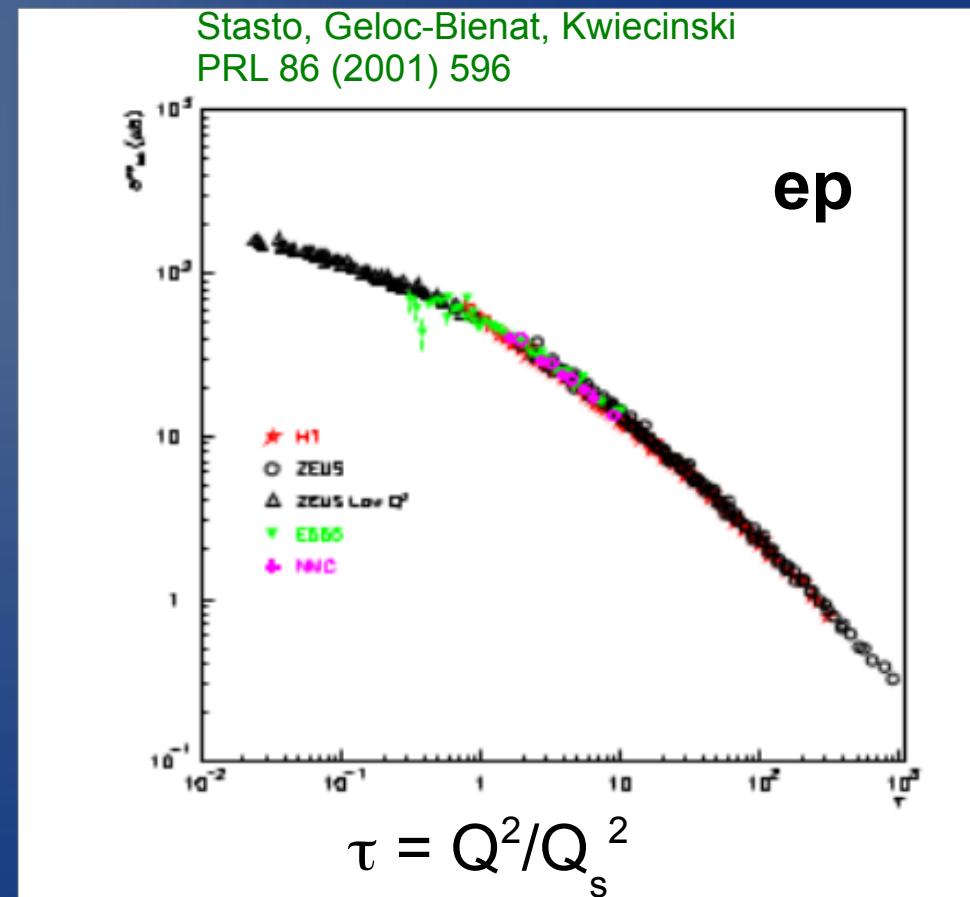
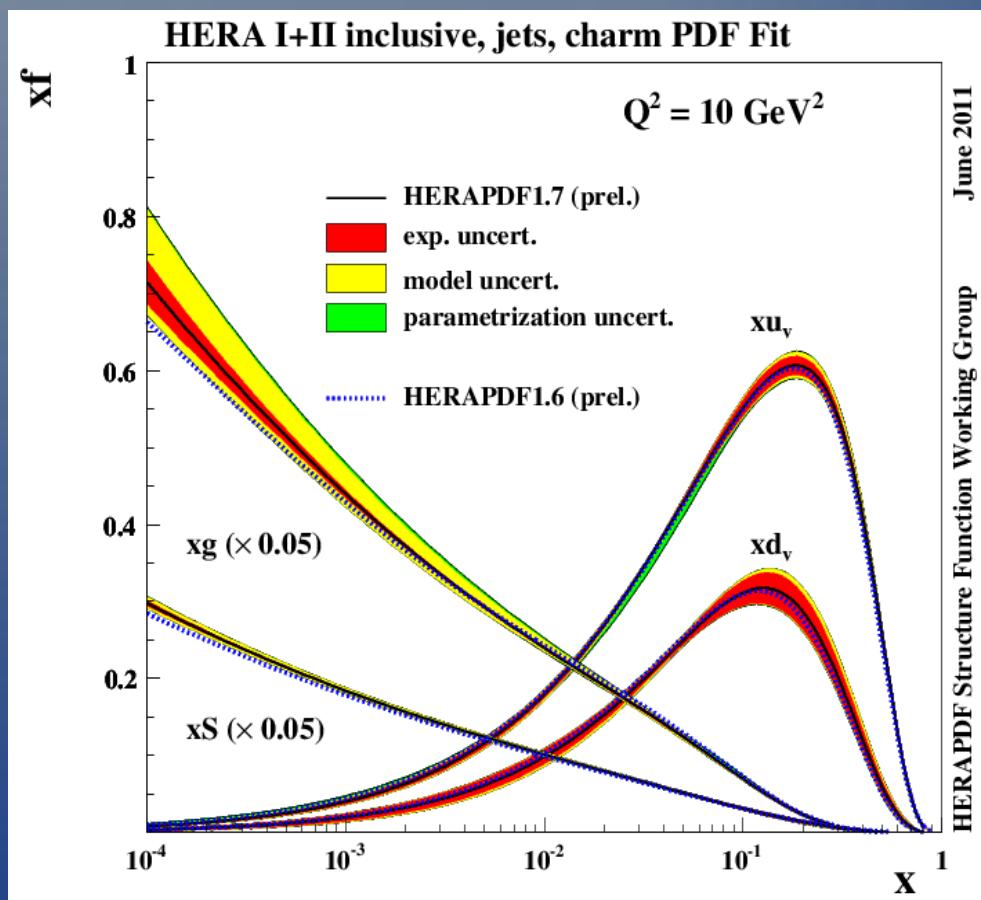


Outline

- Introduction
- Heavy quark production from CGC
- Numerical results
 - J/ψ
 - D meson
 - decay leptons
- Discussion & Summary

Parton saturation

- Dense partonic system at small x is characterized with the new scale, Q_s^2



Parton saturation

- Recombination of partons becomes important at

Phase space density

$$\frac{dN}{d^2rd^2p} \sim \frac{xG(x, Q_s^2)}{\pi R_h^2 Q_s^2(x)} \sim \frac{1}{\alpha_s}$$

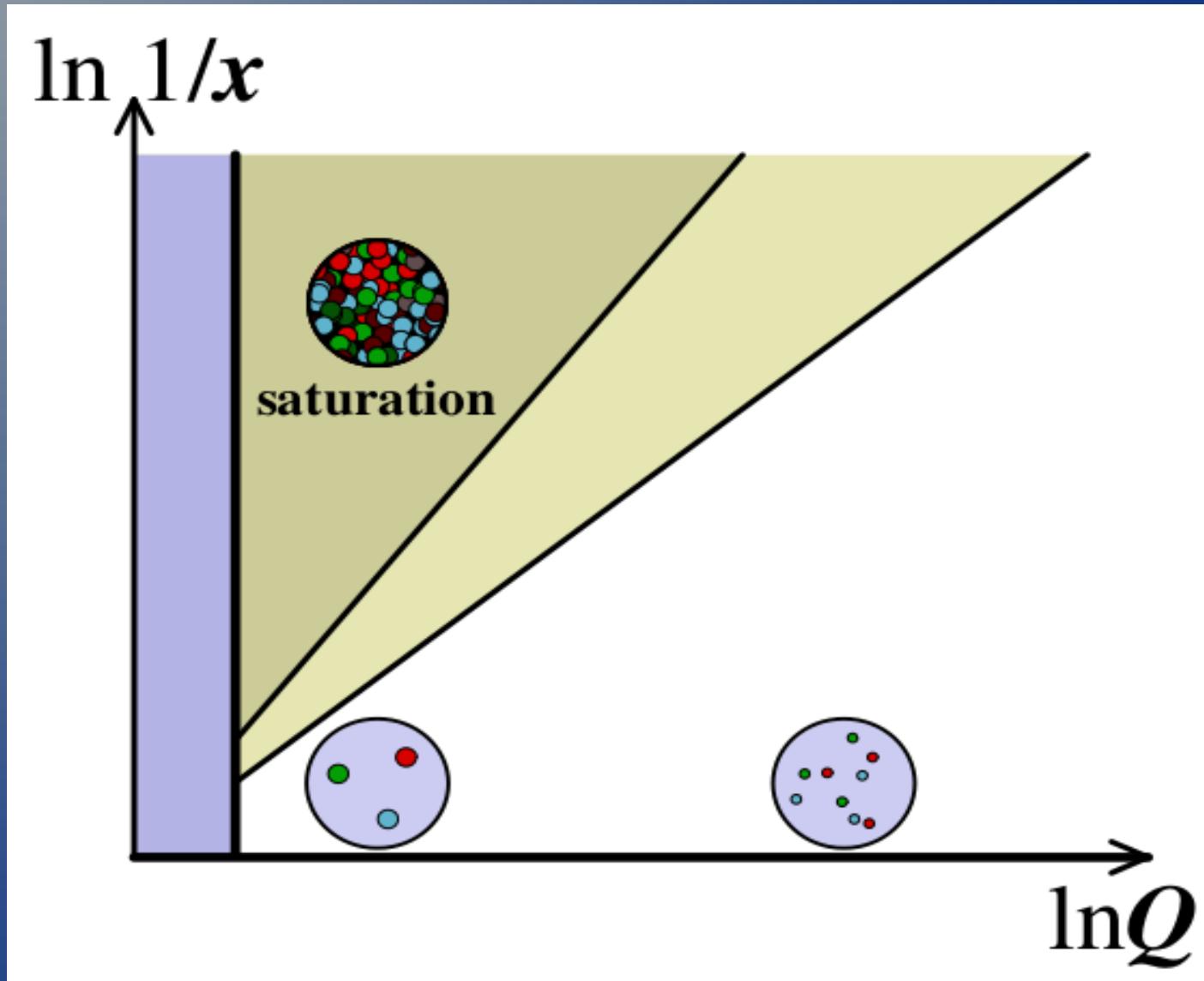
- suggesting a scale, below which saturation operates,

$$Q_s^2(x) \sim \alpha_s \cdot \frac{xG(x, Q_s^2)}{\pi R_h^2}$$

- For a nucleus, atomic number enhances the scale:

$$Q_{s,A}^2(x) \sim \alpha_s \cdot \frac{AxG(x, Q_s^2)}{A^{2/3}\pi R_0^2} = A^{1/3}Q_{s,p}^2$$

Phase diagram



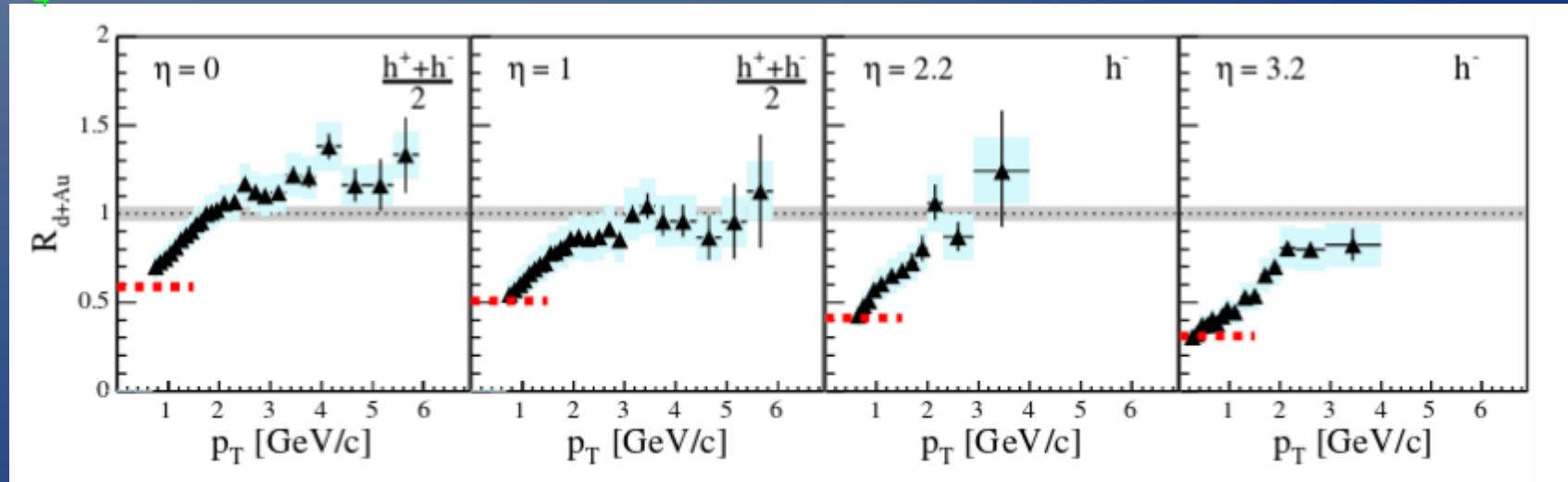
Classic evidence at RHIC

- Saturation is consistent with suppressed particle production in dA, compared to pp, at forward rapidity

BRAHMS,
Phys.Rev.Lett.93:242303,200

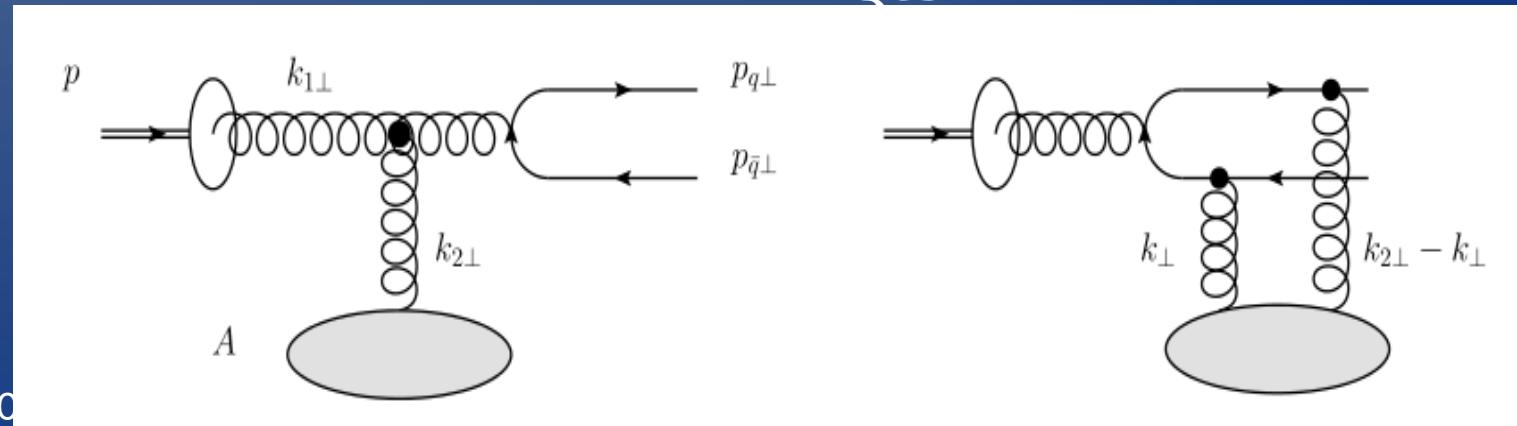
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$$R_{dAu} \equiv \frac{1}{\langle N_{coll} \rangle} \frac{d^2 N^{d+Au} / dp_T d\eta}{d^2 N^{p+p} / dp_T d\eta}$$



Heavy quark production in pA

- Why pA?
 - to study saturation effects, enhanced by big A
- Why heavy quark?
 - produced from gluons
 - can probe saturation Qs ($\sim mQ$)
 - computed in pQCD as $mQ \gg \Lambda_{\text{QCD}}$



Quark pair production in large Nc

- kT-factorization (for small $x_{1,2}$)

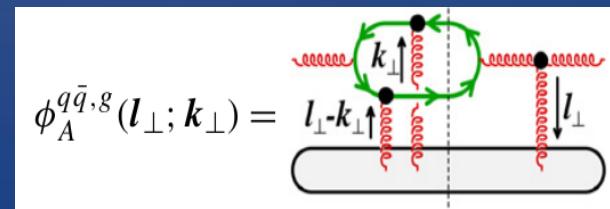
Blaizot-Gelis-Venugopalan 2002
HF-Watanabe 2013

$$\frac{d\sigma_{q\bar{q}}}{d^2 p_{q\perp} d^2 p_{\bar{q}\perp} dy_q dy_{\bar{q}}} = \frac{\alpha_s^2}{16\pi^4 C_F} \int \frac{d^2 k_{2\perp} d^2 k_\perp}{(2\pi)^4} \frac{\Xi(\mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}, \mathbf{k}_\perp)}{k_{1\perp}^2 k_{2\perp}^2} \varphi_{p,x_1}(k_{1\perp}) \phi_{A,x_2}^{q\bar{q},g}(\mathbf{k}_{2\perp}, \mathbf{k}_\perp),$$

- Hybrid formula (for large x_1 and small x_2)

$$\frac{d\sigma_{q\bar{q}}}{d^2 p_{q\perp} d^2 p_{\bar{q}\perp} dy_q dy_{\bar{q}}} = \frac{\alpha_s^2}{16\pi^2 C_F} \int \frac{d^2 k_\perp}{(2\pi)^2} \frac{\Xi_{\text{coll}}(k_{2\perp}, k_\perp)}{k_{2\perp}^2} x_1 G(x_1, \mu) \phi_{A,x_2}^{q\bar{q},g}(k_{2\perp}, k_\perp),$$

- probe the dense gluons through 3pt func.

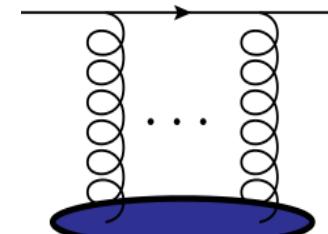


Dipole amplitude

- Wilson line - effective d.o.f.

$$U(z_\perp) = \mathcal{P} \exp \left[ig \int dz^+ A^-(z^+, z) \right]$$

$$\mathcal{N}(\mathbf{x}_\perp, \mathbf{y}_\perp) = \int [D\rho] W_Y[\rho] \left[1 - \frac{1}{N_c} \text{tr}(U(\mathbf{x}_\perp) U^\dagger(\mathbf{y}_\perp)) \right]$$



- BK evolution & I.C. constained by DIS data

$$\frac{\partial \mathcal{N}(r, y)}{\partial y} = \int d^2 \mathbf{r}_1 K^{\text{run}} [\mathcal{N}(r_1, y) + \mathcal{N}(r_2, y) - \mathcal{N}(r, y) - \mathcal{N}(r_1, y) \mathcal{N}(r_2, y)]$$

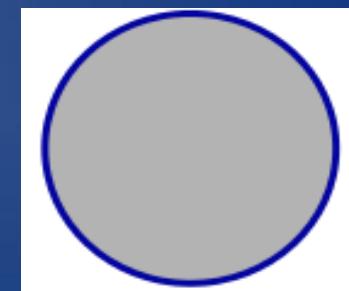
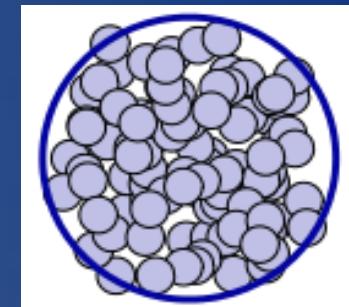
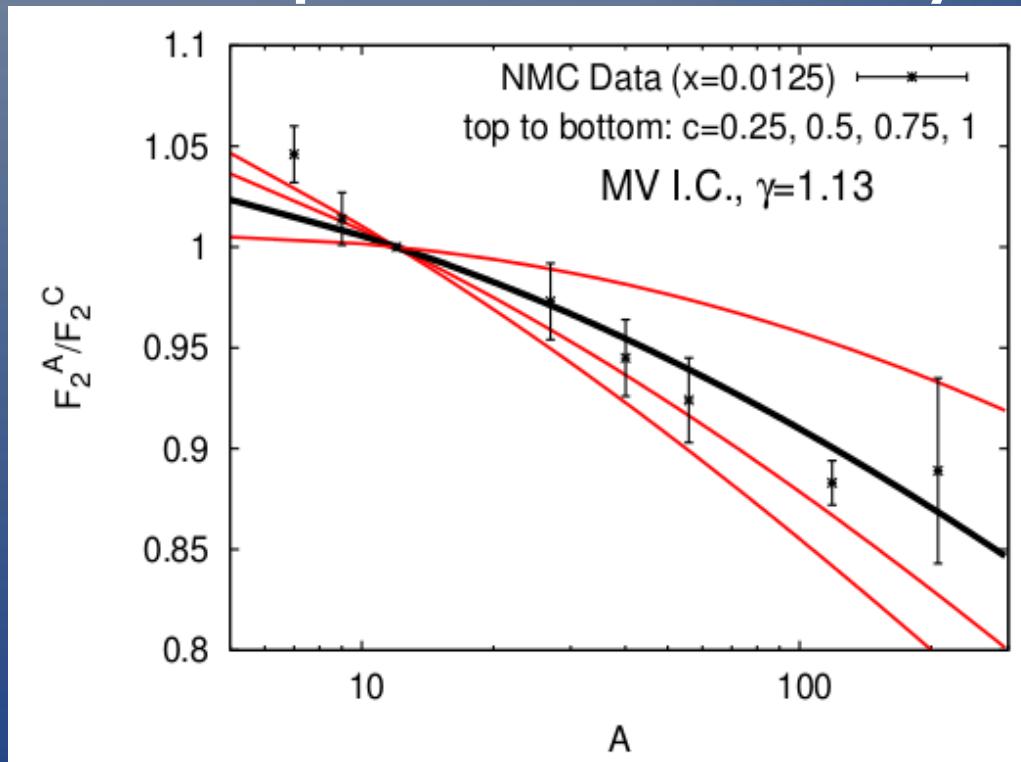
- “gluon dist.” (incl. higher twists):

$$\varphi(k, y) = k^2 \text{F.T.} N(r, y)$$

Choice of Qs for a nucleus

$$Q_{s0A}^2 = c A^{1/3} Q_{s0}^2 = (3?, 6?) Q_{s0}^2$$

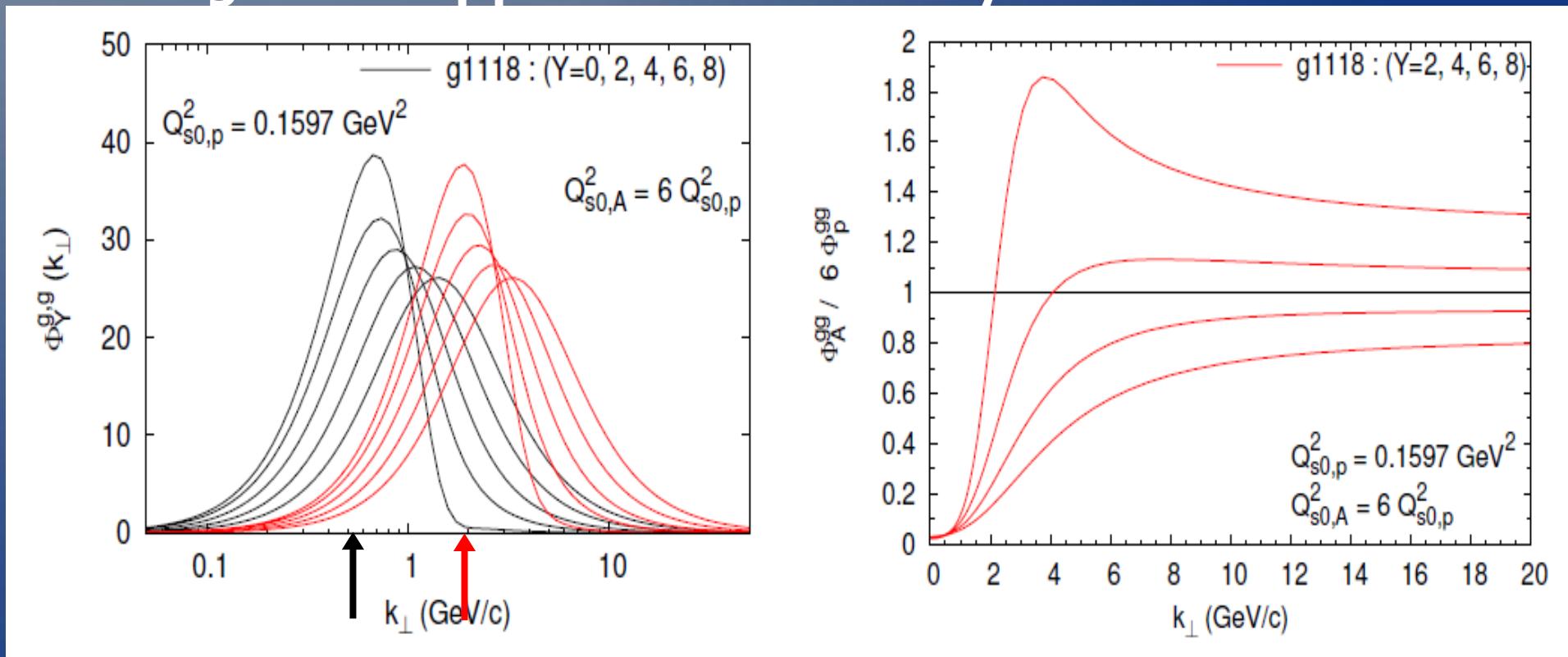
- Effective Qs2 averaged over impact parameter
- c=0.5 is preferred for heavy nuclei



Dusling, et al., NPA836 (2010), 159

Evolution of dipole

- Non-linearity leads to a universal shape of ϕ
- More recombination in dense target
 - gluon suppression in heavy nuclei



Numerical results

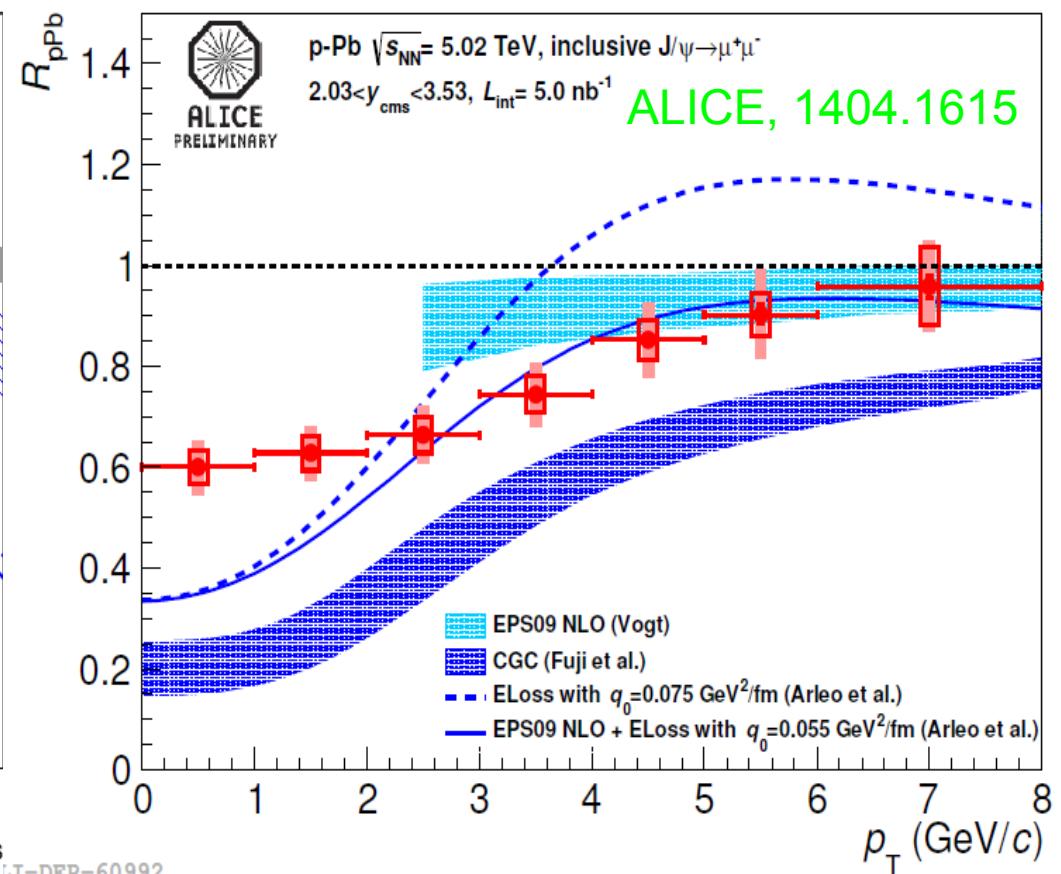
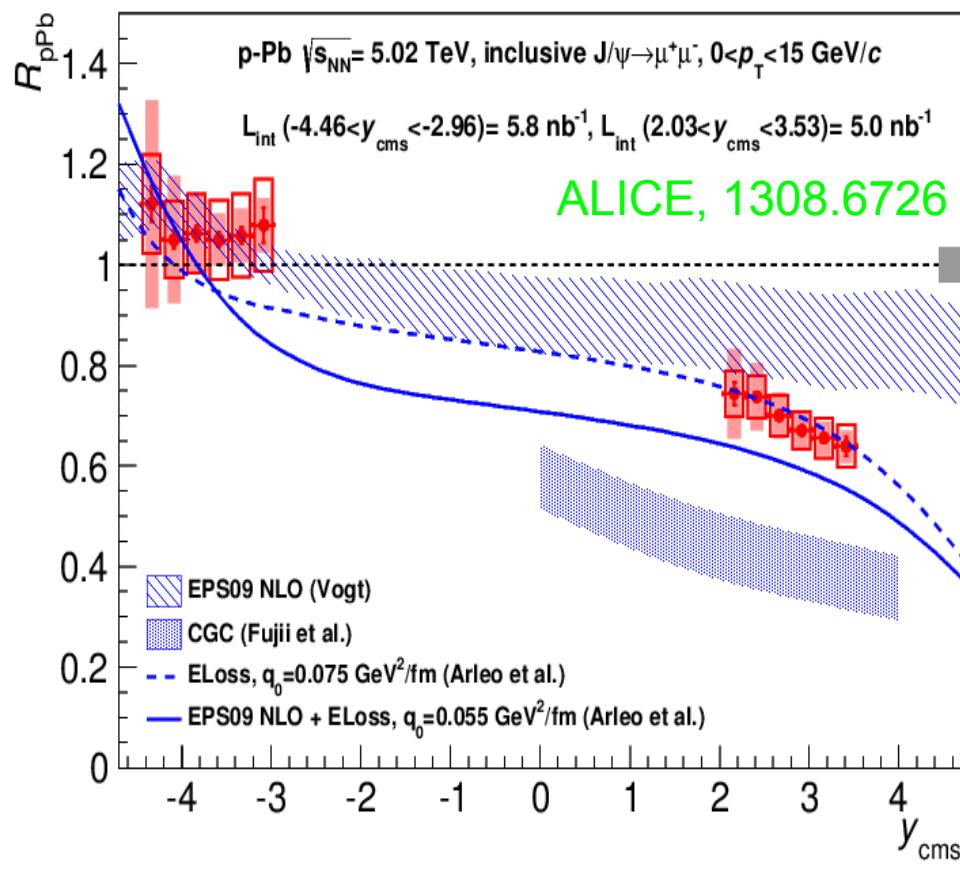
- Quarkonium

RpA for J/ ψ at the LHC

- Color evaporation model –

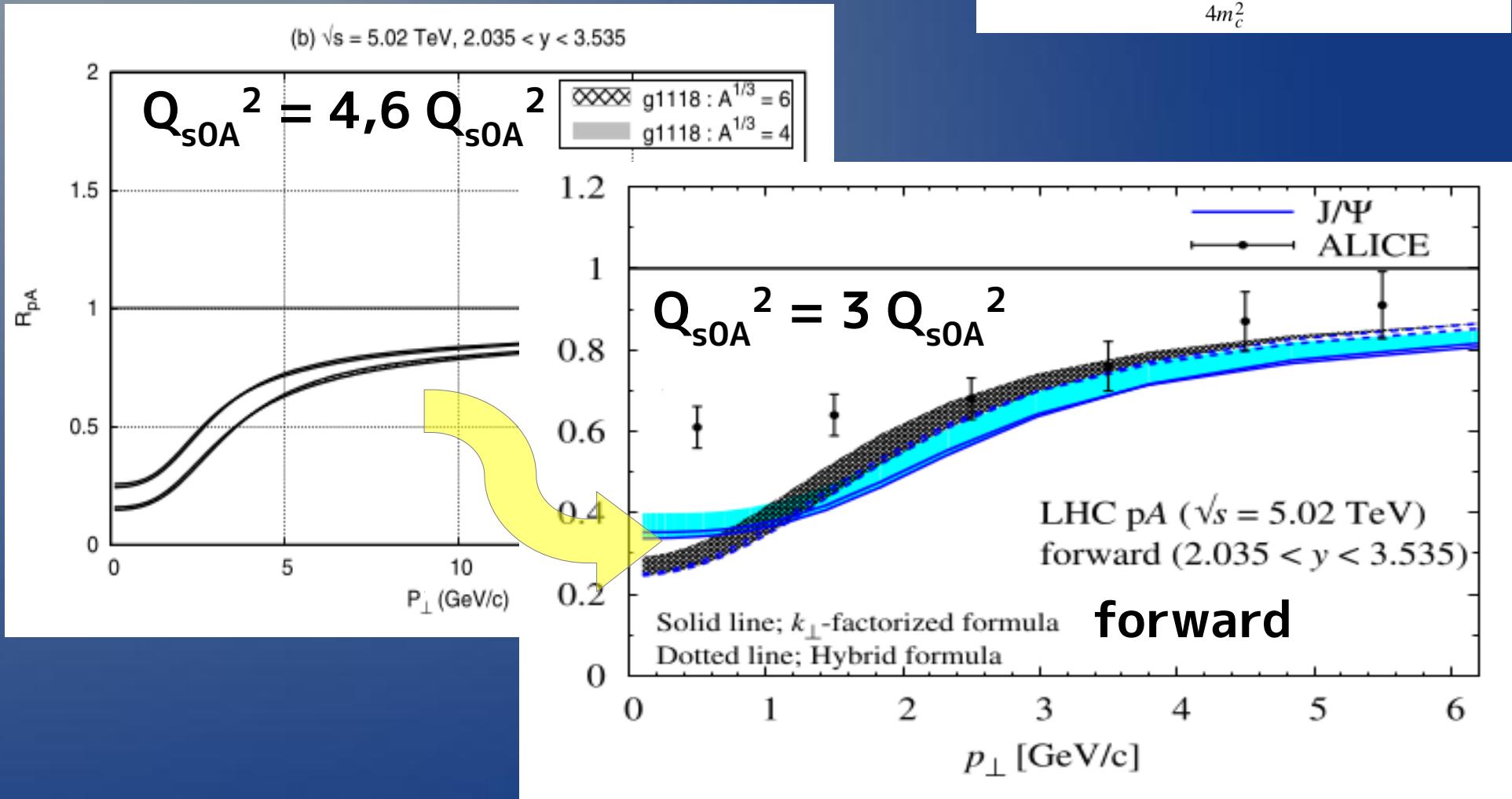
$$Q_{s0A}^2 = (4 - 6) Q_{s0}^2$$

$$\frac{dN_{J/\psi}}{d^2\mathbf{P}_\perp dy} = F_{J/\psi} \int \frac{4M_D^2}{4m_c^2} dM^2 \frac{dN_{c\bar{c}}}{d^2\mathbf{P}_\perp dM^2 dy}$$



RpA for J/ ψ at the LHC

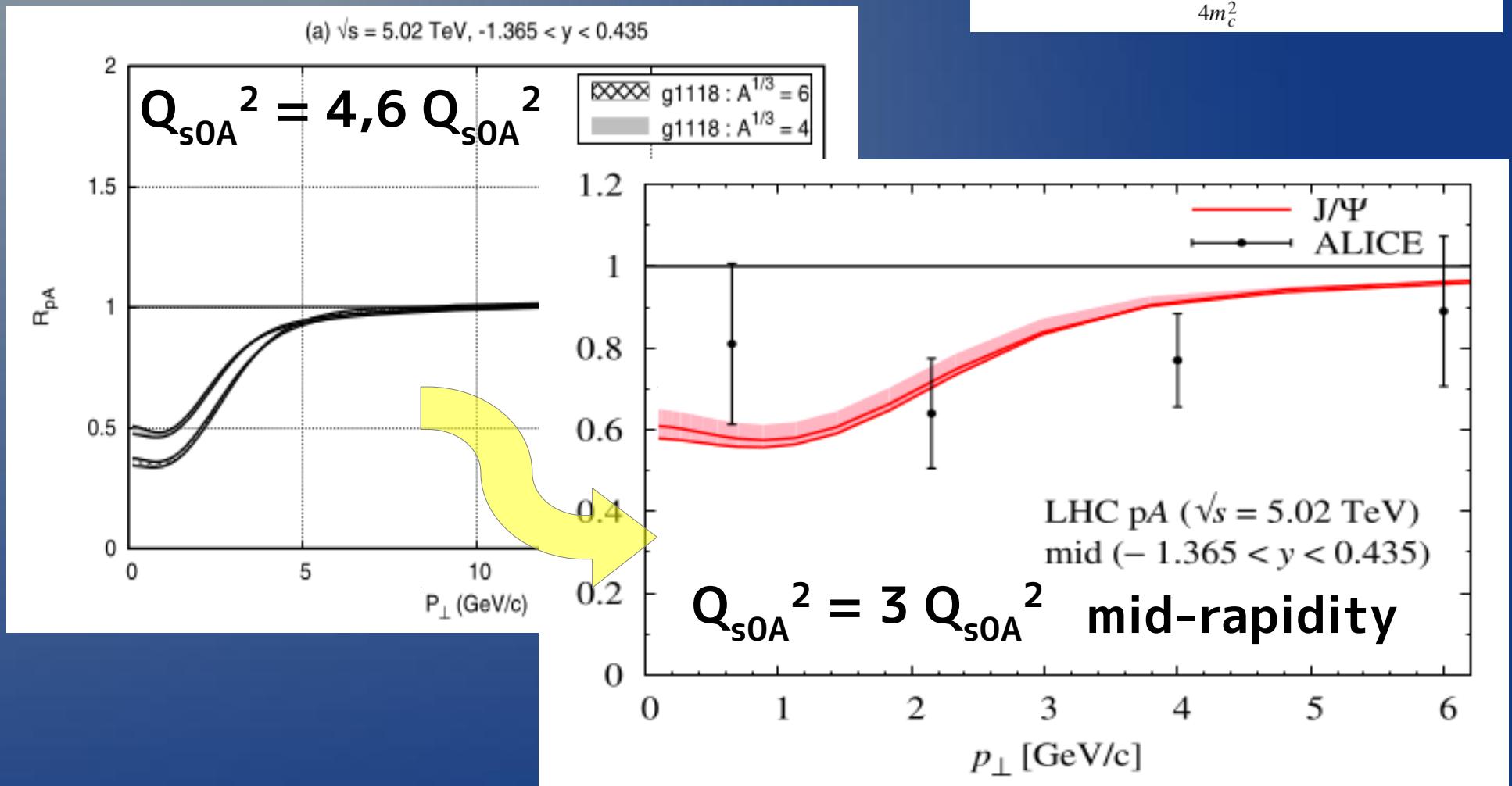
- Color evaporation model -



R_{pA} for J/ ψ at the LHC

- Color evaporation model -

$$\frac{dN_{J/\psi}}{d^2\mathbf{P}_\perp dy} = F_{J/\psi} \int \frac{4M_D^2}{4m_c^2} dM^2 \frac{dN_{c\bar{c}}}{d^2\mathbf{P}_\perp dM^2 dy}$$



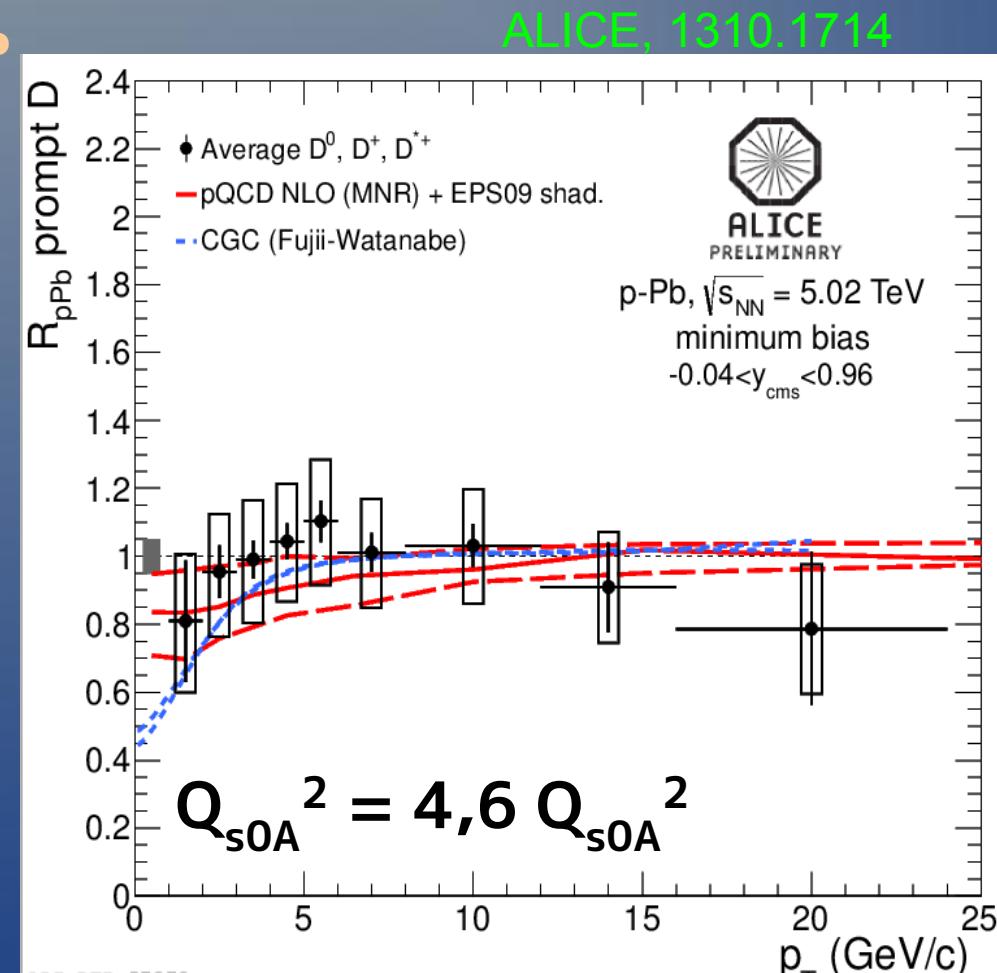
Numerical results

- D meson

R_{pA} for D at the LHC

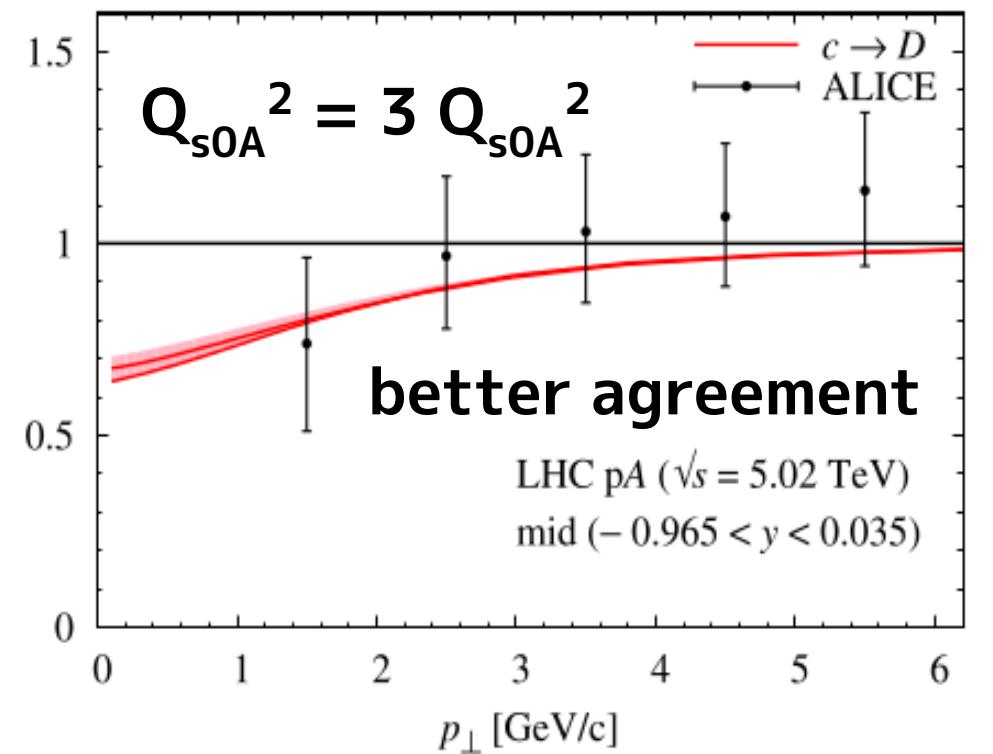
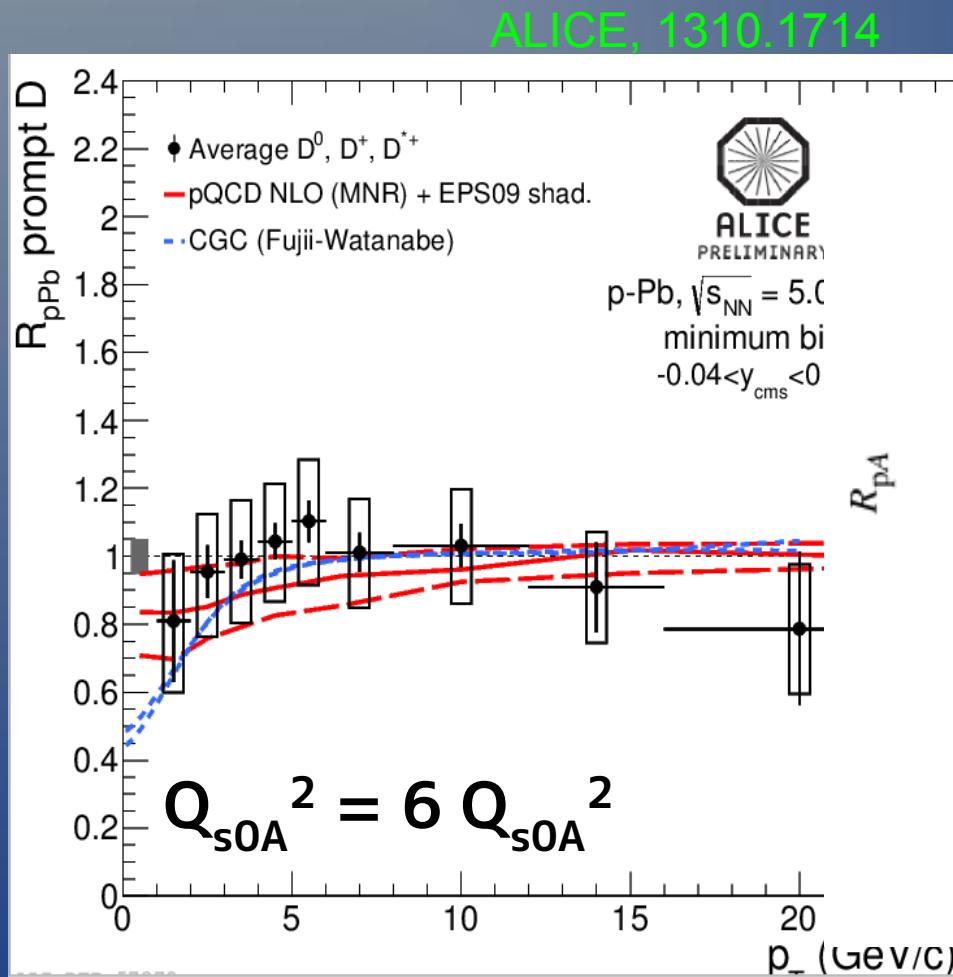
- using frag. fn.

$$\frac{d\sigma_h}{d^2 p_{h\perp} dy} = f_{q \rightarrow h} \int_{z_{\min}}^1 dz \frac{D_q^h(z)}{z^2} \frac{d\sigma_q}{d^2 q_\perp dy}$$



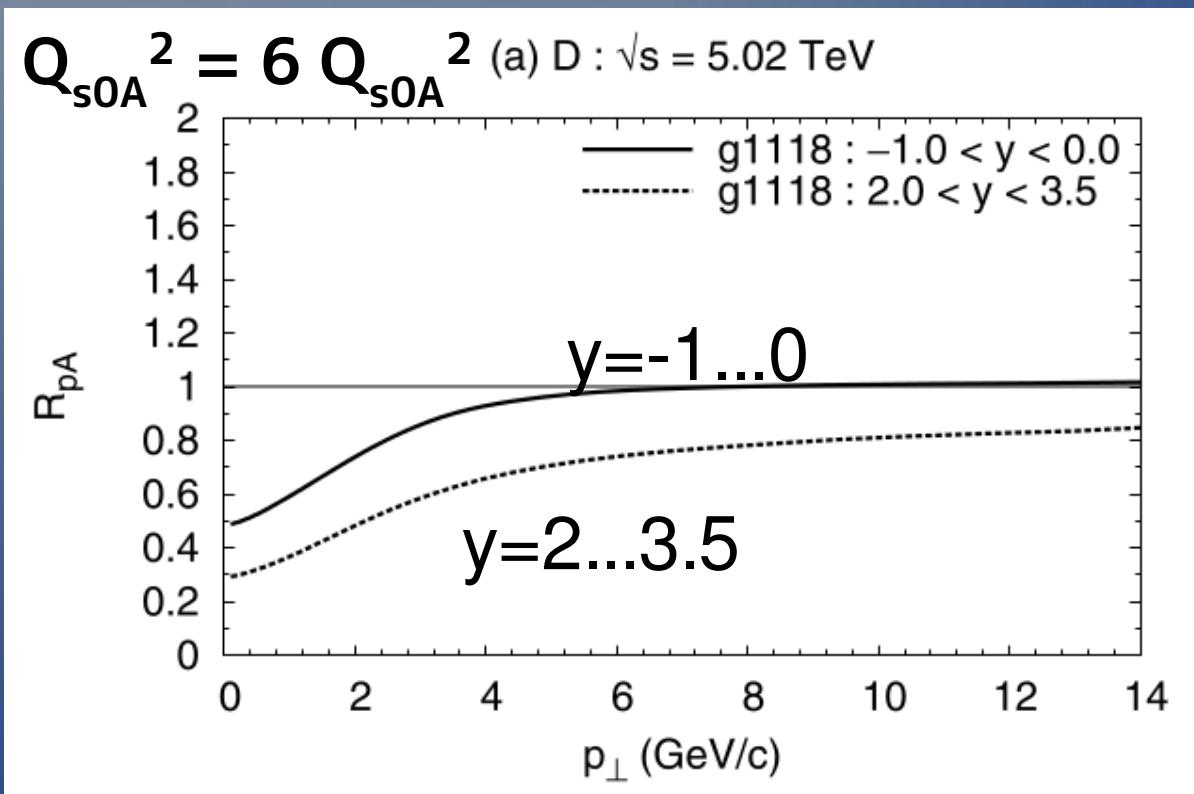
R_{pA} for D at the LHC

- At mid rapidity



R_{pA} for D at the LHC

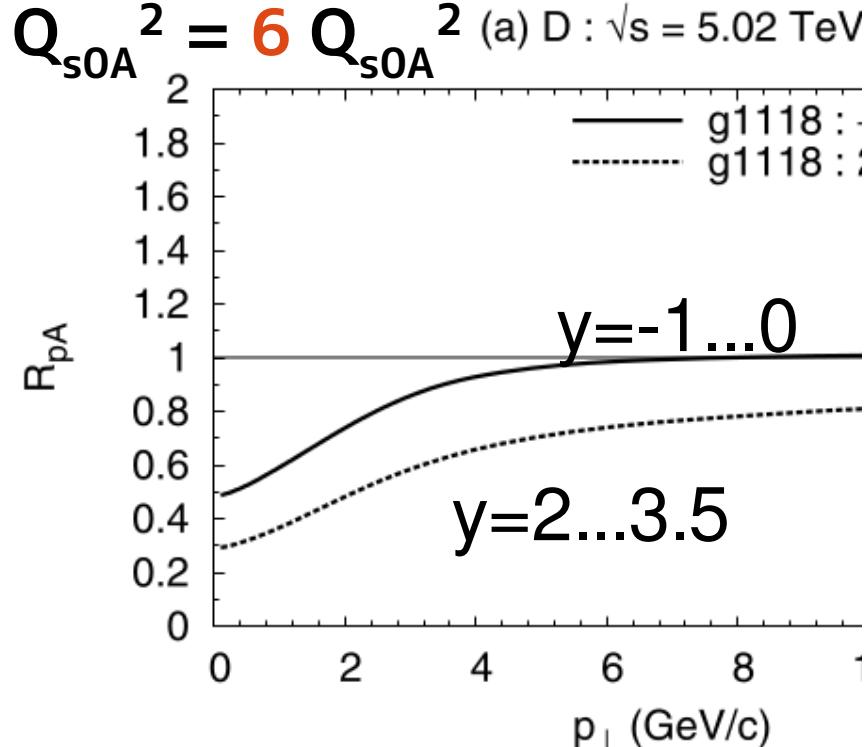
- at forward rapidity



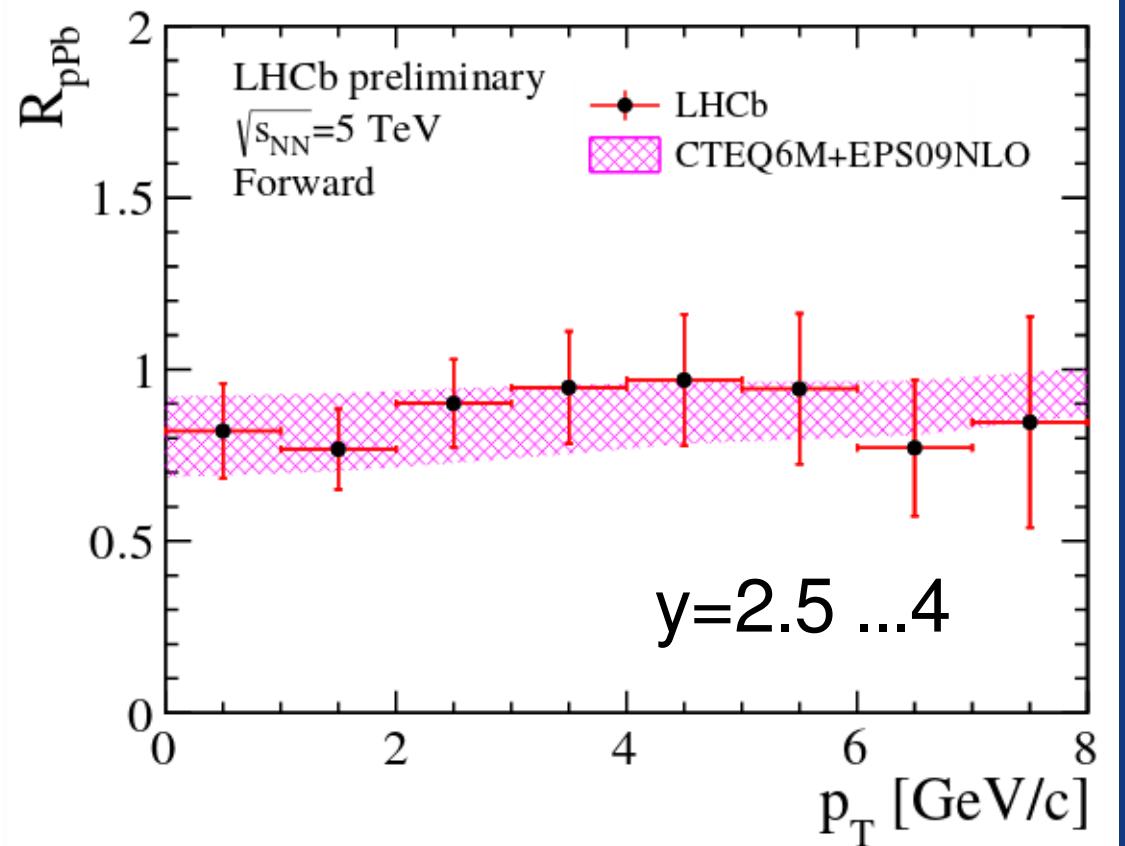
HF, Watanabe, NPA920 (2013)

R_{pA} for D at the LHC

- at forward rapidity



HF, Watanabe, NPA920 (2013)



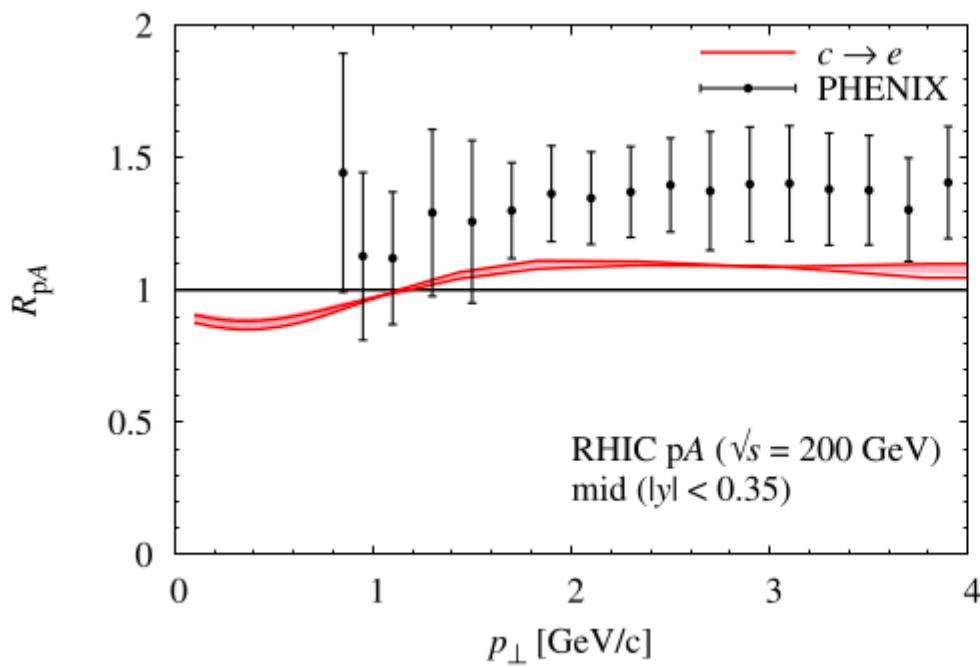
Numerical results

- Decay leptons

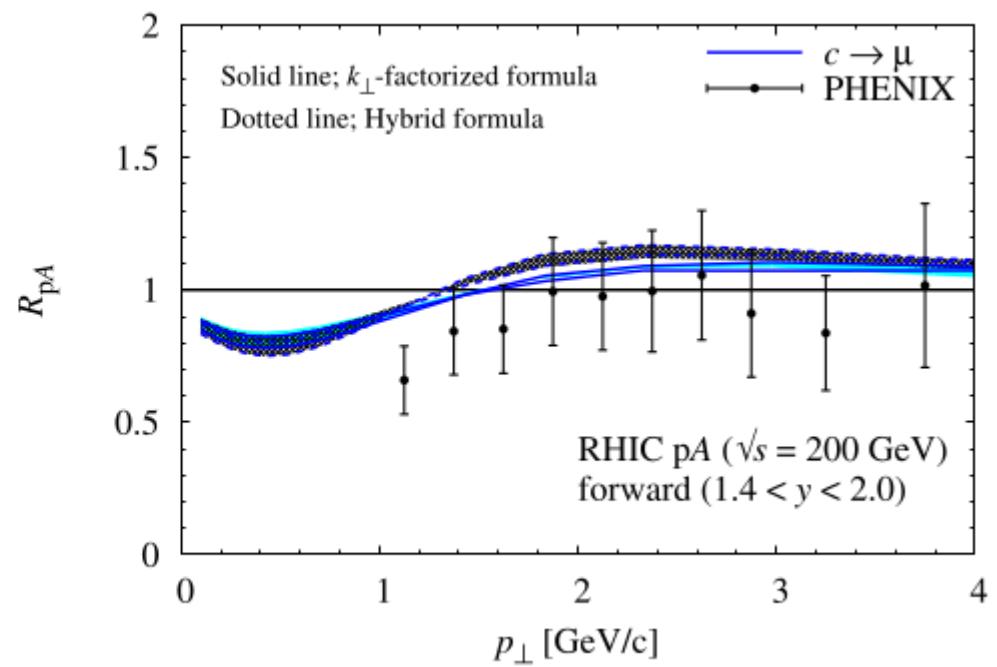
Decay leptons from D at RHIC

- Not inconsistent

RHIC



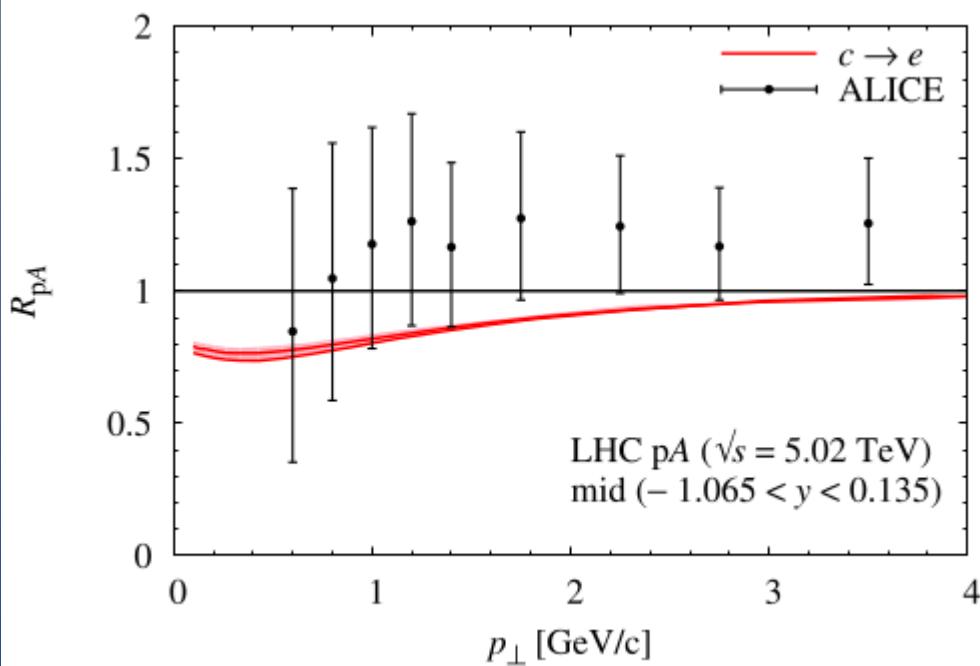
$$Q_{s0A}^2 = 3 Q_{s0A}^2$$



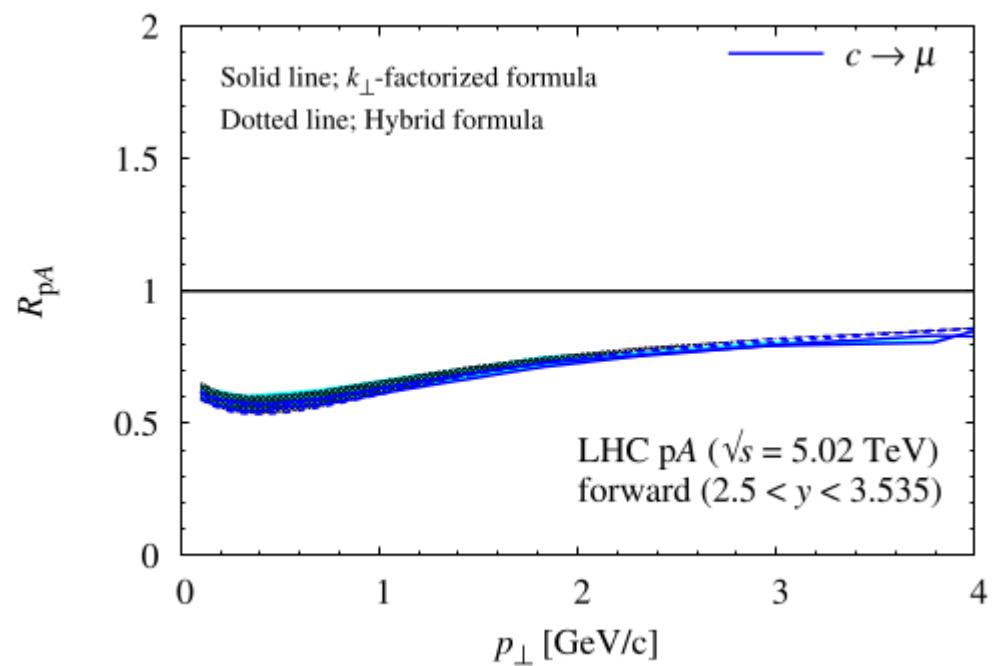
Decay leptons from D at the LHC

- Not inconsistent, but systematically lower than the data

LHC



$$Q_{s0A}^2 = 3 Q_{s0A}^2$$



Discussions

- Results with $Q_{s0A}^2 \sim 3Q_{s0p}^2$ is more consistent to data
- Optical Glauber model [Ducloué et al.](#) in effect corresponds to similar value of Q_{s0A}^2
- Fluctuations in centrality is large in pA
- Should use Monte Carlo to treat a nucleus

Summary

- Updated results with $Q_{s0A}^2 \sim 3Q_{s0p}^2$, which value is consistent with nDIS, describe the data on
 - J/ψ , D mesons, decay leptons
- Monte Carlo description for pA seems important
 - we plan to implement it

Thank you!

- Thank all of you for gathering at Komaba today.
- I hope the discussions today help your research tomorrow.
- Hope to expand the collaboration between France and Japan in our field.
- Look forward!